

Lecture 3 - Rotations and Symmetries

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Active vs. Passive Rotations

Active Rotations

An active rotation is a mapping of the Euclidean space onto itself that keeps at least one point and all distances invariant and preserves orientation.

Passive Rotations

A passive rotation is a coordinate transform from one right handed, orthonormal coordinate system into another one.

Improper Rotations

An improper rotation is a rotation that switches between left handed and right handed coordinate systems.

the matrix $\mathbf{M} = \begin{pmatrix} | & | & | \\ \vec{v}_1 & \vec{v}_2 & \vec{v}_3 \\ | & | & | \end{pmatrix}$

rotates

$$\vec{e}_1 \mapsto \vec{v}_1, \vec{e}_2 \mapsto \vec{v}_2, \vec{e}_3 \mapsto \vec{v}_3$$

\mathbf{M} transforms coordinates from the reference frame $(\vec{v}_1, \vec{v}_2, \vec{v}_3)$ into the reference frame $(\vec{e}_1, \vec{e}_2, \vec{e}_3)$

$\tilde{\mathbf{M}} = -\mathbf{M}$ additionally mirrors all vectors at the origin

Defining Active Rotations

by axis angle

```
v = vector3d(1,0,0)
```

```
w = 90*degree
```

```
r = rotation.byAxisAngle(v,w)
```

```
r = rotation
```

```
Bunge Euler angles in degree
```

```
phi1  Phi  phi2  Inv.
```

```
0     90    0     0
```

Defining Active Rotations

by axis angle

```
v = vector3d(1,0,0)
```

```
w = 90*degree
```

```
r = rotation.byAxisAngle(v,w)
```

by Euler angles

```
r = rotation.byEuler(0,0,pi/2)
```

```
r = rotation
```

```
Bunge Euler angles in degree
```

phi1	Phi	phi2	Inv.
90	0	0	0

Defining Active Rotations

by axis angle

```
v = vector3d(1,0,0)
w = 90*degree
r = rotation.byAxisAngle(v,w)
```

by Euler angles

```
r = rotation.byEuler(0,0,pi/2)
```

by a rotation matrix

```
M = [1 0 0; 0 0 -1; 0 1 0]
r = rotation.byMatrix(M)
```

```
r = rotation
```

Bunge Euler angles in degree

phi1	Phi	phi2	Inv.
0	90	0	0

Defining Active Rotations

by axis angle

```
v = vector3d(1,0,0)
w = 90*degree
r = rotation.byAxisAngle(v,w)
```

by Euler angles

```
r = rotation.byEuler(0,0,pi/2)
```

by a rotation matrix

```
M = [1 0 0; 0 0 -1; 0 1 0]
r = rotation.byMatrix(M)
```

by pairs of vectors

```
u1 = vector3d.Z; v1 = u1
u2 = vector3d.X; v2 = vector3d.Y
r = rotation.map(u1,v1,u2,v2)
```

Defining Active Rotations

by axis angle

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v = vector3d(1,0,0)
w = 90*degree
r = rotation.byAxisAngle(v,w)
```

by Euler angles

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r = rotation.map(u1,v1,u2,v2)
```

```
r = rotation.fit(u,v)
```

Defining Active Rotations

by axis angle

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w = 90*degree
r = rotation.byAxisAngle(v,w)
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by Euler angles

```
r = rotation.byEuler(0,0,pi/2)
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by a rotation matrix

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M = [1 0 0; 0 0 -1; 0 1 0]
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```
u1 = vector3d.Z; v1 = u1
u2 = vector3d.X; v2 = vector3d.Y
r = rotation.map(u1,v1,u2,v2)
```

```
r = rotation.fit(u, v)
```

import rotations

```
r = rotation.load('file.txt',...
  'ColumnNames',...
  {'phi1', 'Phi', 'phi2'})
```

```
r = rotation
size: 1 x 100
```


Defining Active Rotations

by axis angle

```
v = vector3d(1,0,0)
w = 90*degree
r = rotation.byAxisAngle(v,w)
```

by Euler angles

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r = rotation.map(u1,v1,u2,v2)
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r = rotation.fit(u, v)
```

import rotations

```
r = rotation.load('file.txt',...
  'ColumnNames',...
  {'phi1', 'Phi', 'phi2'})
```

random rotations

```
r = rotation.rand(100)
```

```
r = rotation
  size: 1 x 100
```

Defining Active Rotations

by axis angle

```
v = vector3d(1,0,0)
w = 90*degree
r = rotation.byAxisAngle(v,w)
```

by Euler angles

```
r = rotation.byEuler(0,0,pi/2)
```

by a rotation matrix

```
M = [1 0 0; 0 0 -1; 0 1 0]
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```

by pairs of vectors

```
u1 = vector3d.Z; v1 = u1
u2 = vector3d.X; v2 = vector3d.Y
r = rotation.map(u1,v1,u2,v2)
```

```
r = rotation.fit(u, v)
```

import rotations

```
r = rotation.load('file.txt',...
  'ColumnNames',...
  {'phi1', 'Phi', 'phi2'})
```

random rotations

```
r = rotation.rand(100)
```

identity and inversion

```
r = rotation.id
r = rotation.inversion
```

```
r = rotation
```

```
Bunge Euler angles in degree
phi1  Phi  phi2  Inv.
  0    0    0    0
  0    0    0    1
```

Defining Active Rotations

by axis angle

```
v = vector3d(1,0,0)
w = 90*degree
r = rotation.byAxisAngle(v,w)
```

by Euler angles

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by a rotation matrix

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r = rotation.map(u1,v1,u2,v2)
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r = rotation.fit(u, v)
```

import rotations

```
r = rotation.load('file.txt',...
    'ColumnNames',...
    {'phi1', 'Phi', 'phi2'})
```

random rotations

```
r = rotation.rand(100)
```

identity and inversion

```
r = rotation.id
r = rotation.inversion
```

a reflexion

```
r = reflection(vector3d.X)
```

```
r = rotation
```

```
Bunge Euler angles in degree
phi1  Phi  phi2  Inv.
    0  180    0    1
```

Euler Angles

Definition (Euler angles)

Let $\varphi_1, \varphi_2 \in [0, 2\pi]$ and $\Phi \in [0, \pi]$. Then $\varphi_1, \Phi, \varphi_2$ are called the **Euler angles** of the rotation

$$\mathbf{R}(\varphi_1, \Phi, \varphi_2) = \mathbf{R}_{\vec{z}, \varphi_1} \mathbf{R}_{\vec{x}, \Phi} \mathbf{R}_{\vec{z}, \varphi_2}.$$

```
rot = rotation . byEuler (10*degree , 20*degree , 30*degree )
```

- ▶ For every rotation \mathbf{R} there are Euler angles $\varphi_1, \Phi, \varphi_2$ such that $\mathbf{R} = \mathbf{R}(\varphi_1, \Phi, \varphi_2)$.
- ▶ For specific rotations \mathbf{R} the Euler angles are not unique, e.g.

$$\mathbf{R}_{\vec{z}, \omega} = \mathbf{R}(\varphi_1, 0, \omega - \varphi_1)$$

- ▶ Euler angles are the most common way to specify and visualize rotations in texture analysis.
- ▶ The ambiguity makes visualization with respect to Euler angles hard to interpret.

Operations with Rotations

vector rotation

```
v = rot .* u
```

concatenation

```
rot = rot1 .* rot2
```

```
v = rot .* u
```

```
v = rot1 * (rot2 * u)
```

inverse of a rotation / misrotation

```
inv(rot1)
```

```
inv(rot1) .* rot2
```

basic statistics

```
mean(rot)
```

```
mean(rot, 'weights', w)
```

```
std(rot)
```

```
unique(rot)
```

extract Euler angles

```
[phi1, Phi, phi2] = Euler(rot)
```

extract matrix

```
rot.matrix
```

rotation vectors

```
rot.Rodrigues
```

```
rot.homochoric
```

```
rot.cubochoric
```

axis / angle

```
rot.axis, rot.angle
```

```
angle(rot1, rot2)
```

quaternion

```
[a, b, c, d] = double(rot)
```

```
quaternion(rot)
```

Visualizing Rotations

Euler angle space

```
rot = rotation.rand  
plot(rot)
```

axis angle space

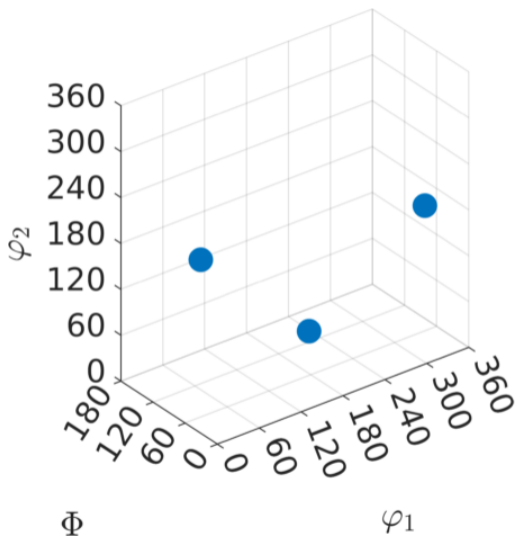
```
plot(rot, 'axisAngle')
```

Euler angle φ_2 sections

```
plotSection(rot, 'phi2')
```

axis angle sections

```
plotSection(rot, 'axisAngle')
```



Visualizing Rotations

Euler angle space

```
rot = rotation.rand  
plot(rot)
```

axis angle space

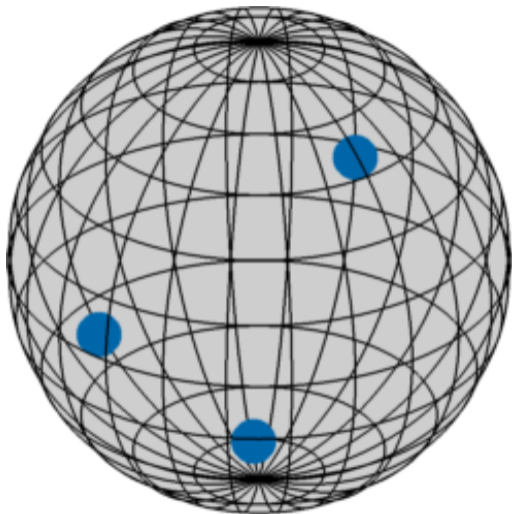
```
plot(rot, 'axisAngle')
```

Euler angle φ_2 sections

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axis angle sections

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Visualizing Rotations

Euler angle space

```
rot = rotation.rand  
plot(rot)
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axis angle space

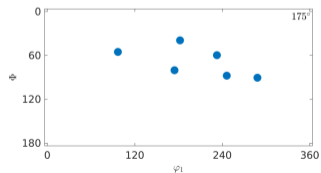
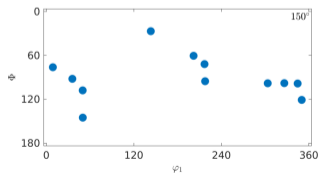
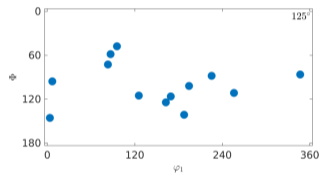
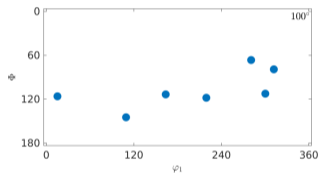
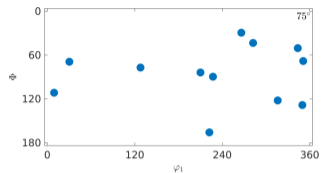
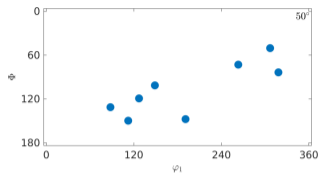
```
plot(rot, 'axisAngle')
```

Euler angle φ_2 sections

```
plotSection(rot, 'phi2')
```

axis angle sections

```
plotSection(rot, 'axisAngle')
```



Visualizing Rotations

Euler angle space

```
rot = rotation.rand  
plot(rot)
```

axis angle space

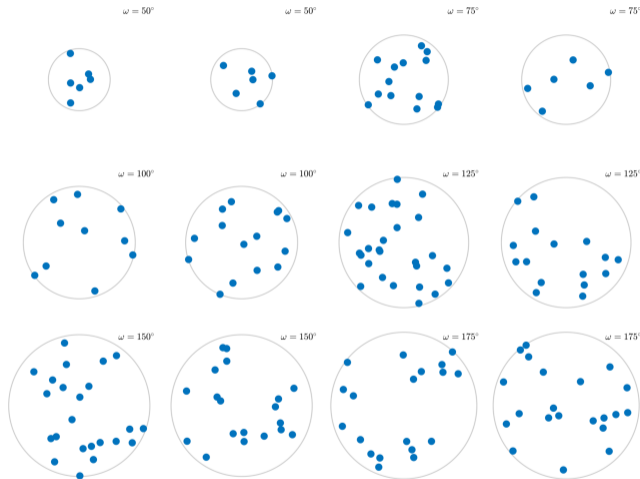
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plot(rot, 'axisAngle')
```

Euler angle φ_2 sections

```
plotSection(rot, 'phi2')
```

axis angle sections

```
plotSection(rot, 'axisAngle')
```



Summary of Rotation Representations

name	notation	space	dimension
matrix	\mathbf{R}	$\mathbb{R}^{3 \times 3}$	9
Euler angles	$(\varphi_1, \Phi, \varphi_2)$	$[0, 2\pi] \times [0, \pi] \times [0, 2\pi]$	3
quaternion	(q_1, q_2, q_3, q_4)	\mathbb{S}^4	4
Rodrigues - Frank	$\tan \frac{\omega}{2} \vec{v}$	\mathbb{R}^3	3
axis angle	$\omega \vec{v}$	\mathbb{B}^3	3
Miller-Bravais Indices	$(hkl)[uvw]$	$\mathbb{S}^2 \times \mathbb{S}^2$	6

Symmetries

Definition

A symmetry is a transformation that keeps something invariant.

Affine Transformations

Definition

Affine transformations are transformations

$$\mathbf{S}(\vec{x}) = \mathbf{R}\vec{x} + \vec{t}$$

that are compositions of a rotation $\mathbf{R} \in O(3)$ and a translation $\vec{t} \in \mathbb{R}^3$.

The set of all affine transformations in the three dimensional space is called **Euclidean motion group** and is denoted by **SE(3)**.

Let $\mathbf{S}_1(\vec{x}) = \mathbf{R}_1\vec{x} + \vec{t}_1$ and $\mathbf{S}_2(\vec{x}) = \mathbf{R}_2\vec{x} + \vec{t}_2$ be two affine transformations. Then their composition

$$\mathbf{S}_2 \circ \mathbf{S}_1(\vec{x}) = \mathbf{S}_2(\mathbf{S}_1(\vec{x})) = \mathbf{R}_2\mathbf{R}_1\vec{x} + \mathbf{R}_2\vec{t}_1 + \vec{t}_2$$

is again an affine transformation.

Affine Transformations

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is again an affine transformation.

Crystal Symmetries

Definition

The subset $\mathcal{E} \subset SE(3)$ of all affine transformations that keep the atom lattice invariant is called **space group** of the crystal.

The space group \mathcal{S} together with the composition \circ is a group, since for any two symmetries $\mathbf{S}_1, \mathbf{S}_2 \in \mathcal{E}$ we have $\mathbf{S}_1 \circ \mathbf{S}_2 \in \mathcal{E}$.

Definition

Let $\mathbf{S}_1, \dots, \mathbf{S}_n \in SE(3)$ arbitrary affine transformations. Then we denote by $\langle \mathbf{S}_1, \dots, \mathbf{S}_n \rangle$ the smallest subgroup of $SE(3)$ that contains the transformations $\mathbf{S}_1, \dots, \mathbf{S}_n$ and call it the **group generated by $\mathbf{S}_1, \dots, \mathbf{S}_n$** .

Example

The group generated by the rotation $\mathbf{R}_{\vec{z}, 120^\circ}$ about 120° about the z-axis is

$$\langle \mathbf{R}_{\vec{z}, 120^\circ} \rangle = \{ \mathbf{R}_{\vec{z}, 120^\circ}, \mathbf{R}_{\vec{z}, 120^\circ}^2, \mathbf{R}_{\vec{z}, 120^\circ}^3 \} = \{ \mathbf{R}_{\vec{z}, 0^\circ}, \mathbf{R}_{\vec{z}, 120^\circ}, \mathbf{R}_{\vec{z}, 240^\circ} \}$$

Crystal Symmetries

Definition

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The space group \mathcal{S} together with the composition \circ is a group, since for any two symmetries $\mathbf{S}_1, \mathbf{S}_2 \in \mathcal{E}$ we have $\mathbf{S}_1 \circ \mathbf{S}_2 \in \mathcal{E}$.

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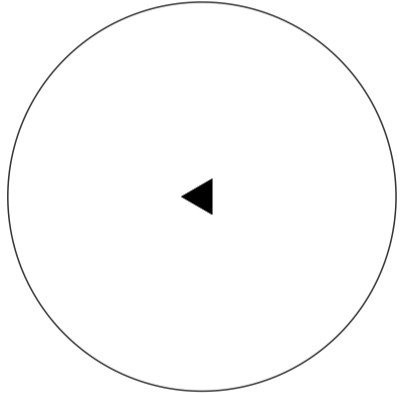
Generate Point Groups in MTEX

```
Z = vector3d.Z  
rot = rotation.byAxisAngle(Z, 120*degree)  
cs = crystalSymmetry.byElements(rot)
```

```
cs.rot = rotation  
size: 3 x 1
```

Bunge Euler angles in degree

phi1	Phi	phi2	Inv.
240	0	0	0
120	0	0	0
0	0	0	0



Generate Point Groups in MTEX

```
Z = vector3d.Z  
rot = rotation.byAxisAngle(Z, 120*degree)  
cs = crystalSymmetry.byElements(rot)
```

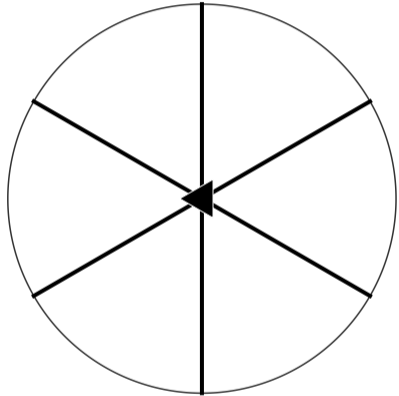
add some mirroring

```
m = reflection(vector3d.X)  
cs = crystalSymmetry.byElements([rot, m])
```

```
cs.rot = rotation  
size: 6 x 1
```

Bunge Euler angles in degree

phi1	Phi	phi2	Inv.
240	180	0	1
0	180	240	1
120	180	120	1
240	0	0	0
120	0	0	0
0	0	0	0



Generate Point Groups in MTEX

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Z = vector3d.Z  
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```

add some mirroring

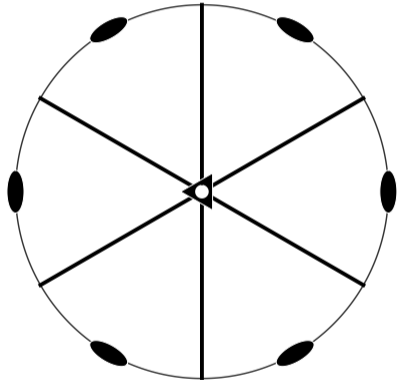
```
m = reflection(vector3d.X)  
cs = crystalSymmetry.byElements([rot, m])
```

add the inversion

```
cs = cs.add(rotation.inversion)
```

```
cs = crystalSymmetry
```

```
symmetry      : -3m1  
elements      : 12  
a, b, c       : 1, 1, 1  
reference frame: X||a, Y||b, Z||c
```



Generate Point Groups in MTEX

```
Z = vector3d.Z  
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```

add some mirroring

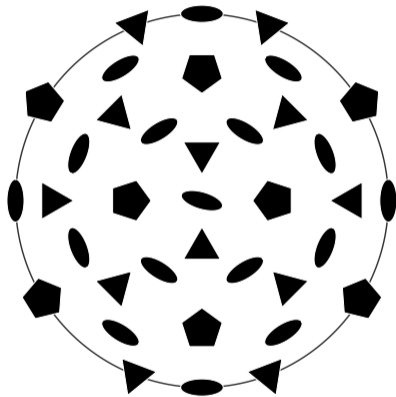
```
m = reflection(vector3d.X)  
cs = crystalSymmetry.byElements([rot, m])
```

add the inversion

```
cs = cs.add(rotation.inversion)
```

some quasi symmetry

```
r2 = rotation.byAxisAngle(Z, 180*degree)  
a5 = vector3d.byPolar(31.7171*degree, 0)  
r5 = rotation.byAxisAngle(a5, 72*degree)  
cs = crystalSymmetry.byElements([r2, r5])
```



Point Groups

Definition

Let $\mathcal{E} = \{(\mathbf{R}_1, \vec{t}_1), (\mathbf{R}_2, \vec{t}_2), \dots\}$ be the space group of a crystal structure. Then $\mathcal{P} = \{\mathbf{R}_1, \mathbf{R}_2, \dots\}$ is a subgroup of $O(3)$ and called the **point group** of the crystal structure.

- ▶ The space group describes the symmetries of an infinite periodic crystal lattice.
- ▶ The point group describes the symmetries of the finite unit cell.

Goal:

Characterize and describe all possible space groups \mathcal{E} and all possible point groups \mathcal{P} .

Result:

There are exactly 230 different space groups and 32 different point groups.

Point Groups

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Result:

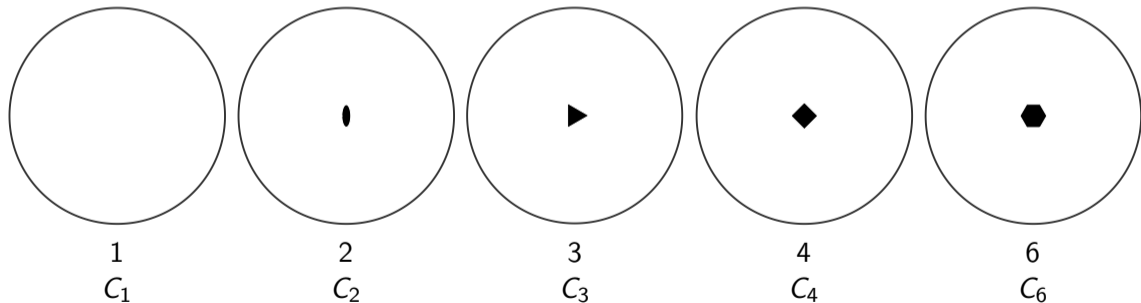
There are exactly 230 different space groups and 32 different point groups.

Symmetry Operations

name	affine transformation	Hermann–Mauguin symbol	graphical symbol
rotational axis	$\mathbf{R}_{\vec{d}, \frac{360^\circ}{n}}$	2, 3, 4, 6	
inversion	$-\mathbf{I}$	$\bar{1}$	
mirror plane	$-\mathbf{R}_{\vec{n}, 180}$	$m = \bar{2}$	
rotoinversion axis	$-\mathbf{R}_{\vec{d}, \frac{360^\circ}{n}}$	$\bar{3}, \bar{4}, \bar{6}$	
translation	\vec{t}		
screw axis	$\mathbf{R}_{\vec{d}, \frac{360^\circ}{n}} + \vec{d}$	$2_1, 3_1, 3_2, 4_1, 4_2, 4_3$ $6_1, 6_2, 6_3, 6_4, 6_5$	
glide plane	$-\mathbf{R}_{\vec{d}, \frac{180^\circ}{n}} + \vec{d}$	a, b, c, n, d	

Cyclic Groups

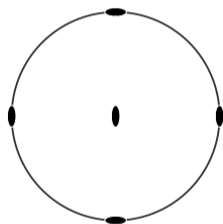
- ▶ $\mathbf{S} \in \mathcal{P} \implies \mathbf{S}^n \in \mathcal{P}$, for all $n \in \mathbb{Z}$
- ▶ $\mathbf{R}_{\vec{d}, \frac{1}{m}360^\circ} \in \mathcal{P} \implies \mathbf{R}_{\vec{d}, \frac{n}{m}360^\circ} \in \mathcal{P}$, $n = 0, \dots, m - 1$.
- ▶ Only rotational axes of order 2, 3, 4, 6 are compatible with periodic lattices



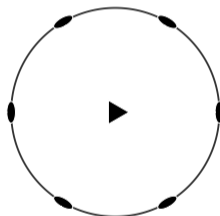
Dieder Groups

Two two fold symmetry axis \vec{a} , \vec{b} at an angle generate a perpendicular symmetry axis

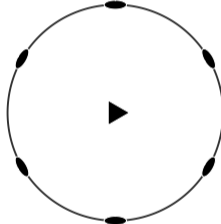
► $\mathbf{R}_{\vec{a},\pi}, \mathbf{R}_{\vec{b},\pi} \in \mathcal{P}, \angle(\vec{a}, \vec{b}) = \frac{\pi}{n} \implies \mathbf{R}_{\vec{a} \times \vec{b}, \frac{2\pi}{n}} \in \mathcal{P}$



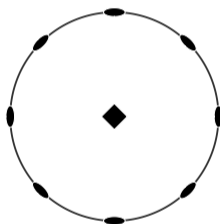
222
 D_2



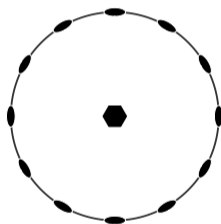
321
 D_3



312
 D_3



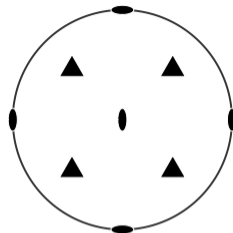
422
 D_4



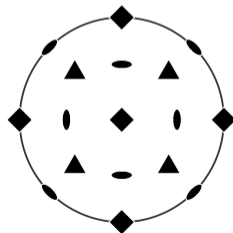
622
 D_6

Tetragonal and Cubic Symmetry

- ▶ \vec{a} - m -fold symmetry axis, $m > 2$
- ▶ \vec{b} - n -fold symmetry axis
- ▶ $\implies \mathbf{R}_{\vec{a}, \frac{k}{m}2\pi} \vec{b}$, $k = 1, \dots, m$ are n -fold symmetry axes
- ▶ $\implies \mathbf{R}_{\vec{b}, \frac{k}{n}2\pi} \vec{a}$, $k = 1, \dots, n$ are m -fold symmetry axes
- ▶ $\implies \mathbf{R}_{\vec{b}, \frac{2k+1}{2n}2\pi} \vec{a}$, $k = 0, \dots, n$ are 2-fold symmetry axes
- ▶ assume \vec{a} , \vec{b} have minimum angle in \mathcal{P}
- ▶ $\implies \mathbf{R}_{\vec{a}, \frac{k}{m}2\pi} \vec{b}$, $k = 1, \dots, m$ form a regular spherical polygon $P_{\vec{a}}$
- ▶ applying all symmetry operations of \mathcal{P} to $P_{\vec{a}}$ will cover the whole sphere by disjoint copies of $P_{\vec{a}}$
- ▶ \implies the copies of $P_{\vec{a}}$ form a Platonic solid
- ▶ \implies only tetrahedron and cube (octahedron) are relevant

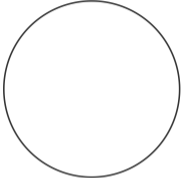
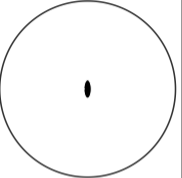
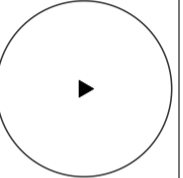
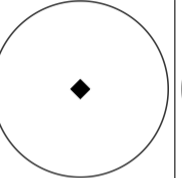
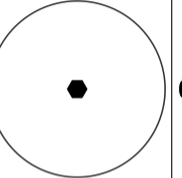
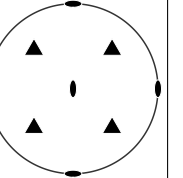
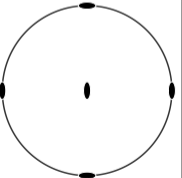
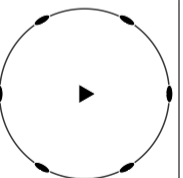
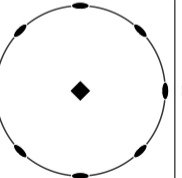
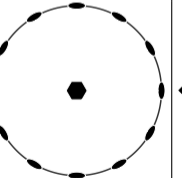
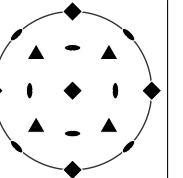


23 (T)

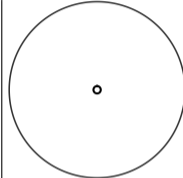
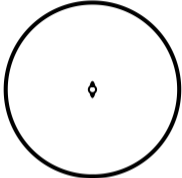
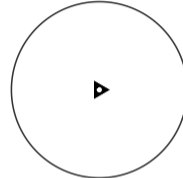
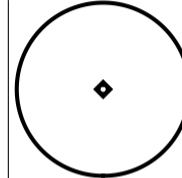
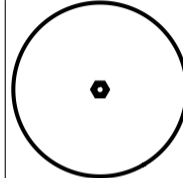
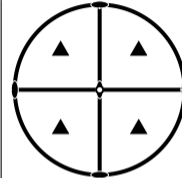
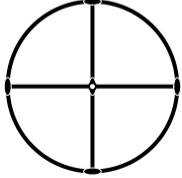
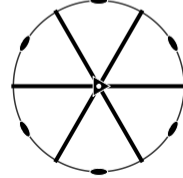
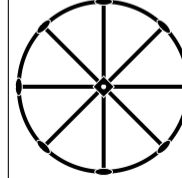
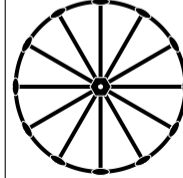
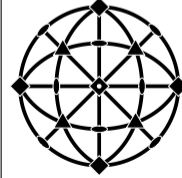


432 (O)

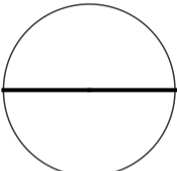
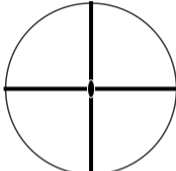
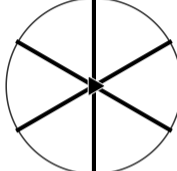
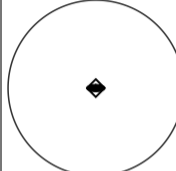
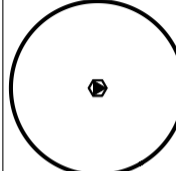
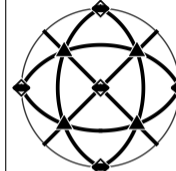
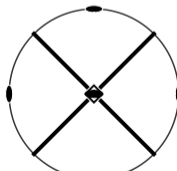
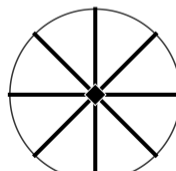
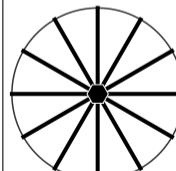
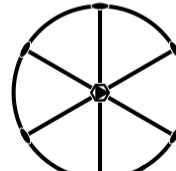
All 11 Enantiomorphic Symmetry Groups

triclinic	monoclinic	trigonal	tetragonal	hexagonal	cubic
 1 (C_1)	 2 (C_2)	 3 (C_3)	 4 (C_4)	 6 (C_6)	 23 (T)
	orthorhombic				
	 222 (D_2)	 321 (D_3)	 422 (D_4)	 622 (D_6)	 432 (O)

All 11 Laue Symmetry Groups

<p>triclinic</p>  <p>$\bar{1}$ (C_i)</p>	<p>monoclinic</p>  <p>$\frac{2}{m}$ (C_{2h})</p>	<p>trigonal</p>  <p>$\bar{3}$ (C_{3i})</p>	<p>tetragonal</p>  <p>$\frac{4}{m}$ (C_{4h})</p>	<p>hexagonal</p>  <p>$\frac{6}{m}$ (C_{6h})</p>	<p>cubic</p>  <p>$m\bar{3} = \frac{2}{m}\bar{3}$ (T_h)</p>
	<p>orthorhombic</p>  <p>mmm (D_{2h})</p>	 <p>$\bar{3}m = \bar{3}\frac{2}{m}$ (D_{3d})</p>	 <p>$4/mmm$ (D_{2h})</p>	 <p>$6/mmm$ (D_{6h})</p>	 <p>$m\bar{3}m$ (O_h)</p>

10 mixed groups

monoclinic	orthorhombic	trigonal	tetragonal	hexagonal	cubic
 <p>$\bar{2} = m$ (C_s)</p>	 <p>$mm2$ (C_{2v})</p>	 <p>$3m$ (C_{3v})</p>	 <p>$\bar{4}$ (S_4)</p>	 <p>$\bar{6}$ (C_{3h})</p>	 <p>$\bar{4}3m$ (T_d)</p>
		 <p>$\bar{4}2m$ (D_{2d})</p>	 <p>$4mm$ (C_{4v})</p>	 <p>$6mm$ (C_{6v})</p>	 <p>$\bar{6}m2$ (D_{3h})</p>

Summary Symmetry Groups

- ▶ 11 purely rotational (enantiomorphic) groups:
 - used in most software,
 - only proper rotations are considered
- ▶ 11 Laue groups (with inversion center, centrosymmetric groups):
 - correct models for most diffraction experiments, e.g. X-ray diffraction, EBSD
 - only few physical properties are not centrosymmetric e.g., piezoelectricity
- ▶ 10 mixed groups
 - without inversion center
 - each with equally many proper and improper symmetry elements
- ▶ 32 different point groups
 - do **not** represent actual symmetries of the atom lattice - only modulu translation
- ▶ 230 different point groups
 - The translation vectors always coincide with the screw axes.
 - completely described in *International Tables for Crystallography*, 2016

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Defining Crystal Symmetries in MTEX

```
cs = crystalSymmetry('m-3m') % by international symbol  
cs = crystalSymmetry('Oh') % by Schoenflies notation  
cs = crystalSymmetry('Fm-3m') % by space group
```

```
% import CIF file
```

```
cs = crystalSymmetry.load('quartz.cif')
```

```
% download from Crystallography Open Database
```

```
cs = crystalSymmetry.load('5000036')
```

```
cs.properGroup
```

```
cs.Laue
```

```
cs.rot
```

```
plotb2east
```

```
plot(cs)
```

The 14 Bravais lattices

Question: Which lattices are compatible with the symmetry groups?

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Compatibility with the points groups leads to the **7 crystal systems**:

- triclinic
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- tetragonal
- hexagonal
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Compatibility with the space groups leads to the **the 14 Bravais lattices** which additionally separates how the atoms are aligned with the lattice points:

- ▶ primitive
- ▶ base centered
- ▶ body centered
- ▶ face centered

system

P (primitive)

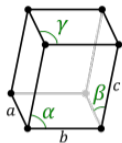
C (base centered)

I (body centered)

F (face centered)

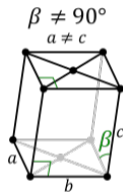
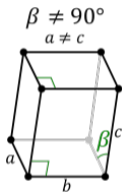
triclinic

$$a \neq b \neq c,$$
$$\alpha \neq \beta \neq \gamma$$



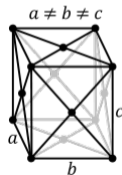
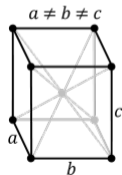
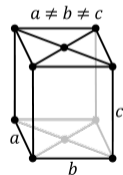
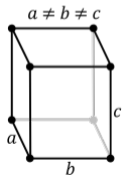
monoclinic

$$a \neq b \neq c,$$
$$\alpha = \gamma = 90^\circ$$
$$\beta > 90^\circ$$



orthorhombic

$$a \neq b \neq c,$$
$$\alpha = \beta = 90^\circ$$
$$\gamma = 90^\circ$$

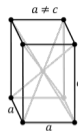
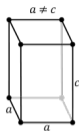


system**P (primitive)****C (base centered)****I (body centered)****F (face centered)****tetragonal**

$$a = b \neq c,$$

$$\alpha = \beta = 90^\circ$$

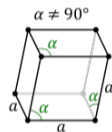
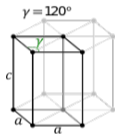
$$\gamma = 90^\circ$$

**trigonal**

$$a = b \neq c,$$

$$\alpha = \beta = 90^\circ$$

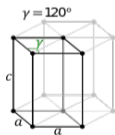
$$\gamma = 120^\circ$$

**hexagonal**

$$a = b \neq c,$$

$$\alpha = \beta = 90^\circ$$

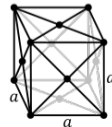
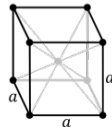
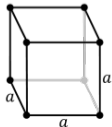
$$\gamma = 120^\circ$$

**cubic**

$$a = b = c,$$

$$\alpha = \beta = 90^\circ$$

$$\gamma = 90^\circ$$



A Practical Example

- ▶ checkout `https://materialsproject.org/materials/mp-2657/#`
- ▶ download cif file

The Ambiguity of the Crystal Coordinate System

the axes of the crystal coordinate system \vec{a} , \vec{b} , \vec{c} are always chosen such that

- ▶ the translations $\mathbf{T}_{\vec{a}}$, $\mathbf{T}_{\vec{b}}$, $\mathbf{T}_{\vec{c}}$ are symmetry elements of the space group
- ▶ \vec{c} is the axis of highest symmetry (except for monoclinic and 23)
- ▶ \vec{a} or \vec{b} are aligned with the symmetry axes perpendicular to \vec{c}

critical symmetries

monoclinic alignment of the two fold axis: 211 , 121 , 112 , $m11$, $1m1$, $11m$,

orthorhombic alignment of the two fold axis: $2mm$, $m2m$, $mm2$

trigonal alignment of the two fold axis: 321 , 312

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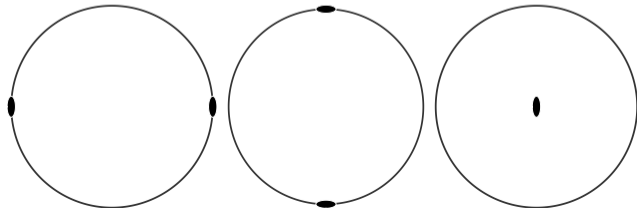
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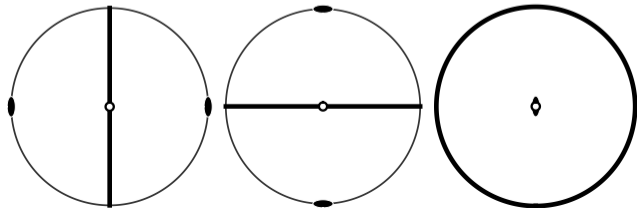
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orthorhombic alignment of the two fold axis: $2mm$, $m2m$, $mm2$

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The Ambiguity of the Crystal Coordinate System

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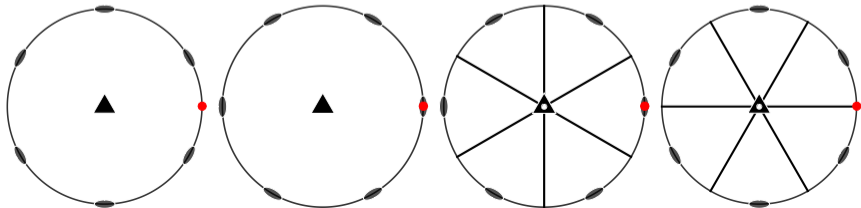
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Symmetrically Equivalent Lattice Directions

The crystal axes \vec{a} , \vec{b} , \vec{c} can be well defined modulo actions of the symmetry group \mathcal{S} .

The axes $(\vec{a}, \vec{b}, \vec{c})$ and $(\mathbf{S}\vec{a}, \mathbf{S}\vec{b}, \mathbf{S}\vec{c})$ are physically indistinguishable for all $\mathbf{S} \in \mathcal{S}$

Two lattice directions

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- ▶ $\langle uvw \rangle$ may contain at maximum $|\mathcal{S}|$ different lattice directions
- ▶ rotational axes have fewer symmetrically equivalent crystal directions, e.g. $\langle 001 \rangle = [001]$
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Symmetrically Equivalent Lattice Planes

For most symmetries the symmetrically equivalent directions comes as permutations of the Miller indices

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Analogously the class of symmetrically equivalent lattice planes $\{hkl\}$ is defined as the set of all lattice planes $(h_2k_2l_2)$ such that there is a symmetry operation $S \in \mathcal{S}$ with

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reference system	single	symmetrically equivalent
lattice point	$\cdot uvw \cdot$	$: uvw :$
lattice direction	$[uvw]$	$\langle uvw \rangle$
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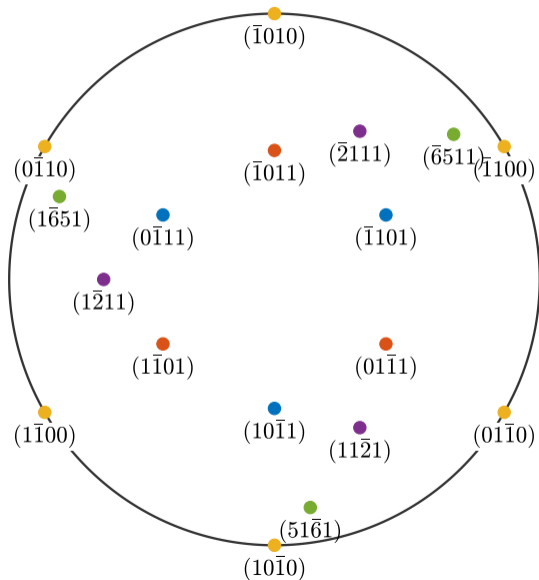
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Solution: 4 digit Miller indices

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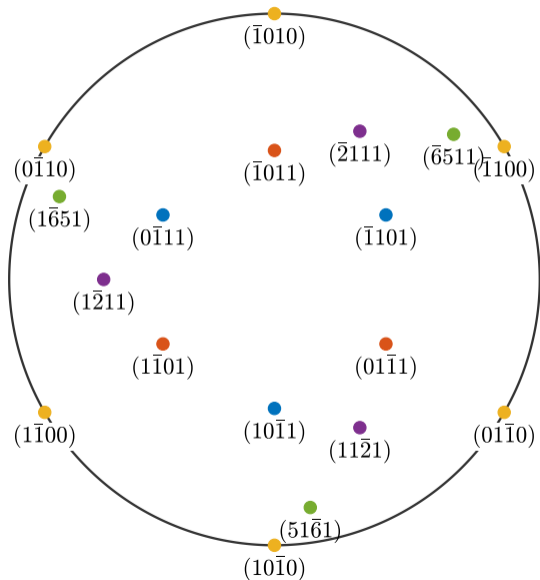
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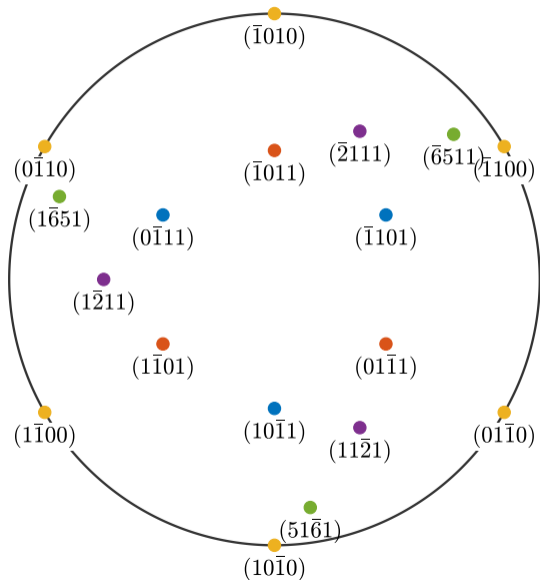
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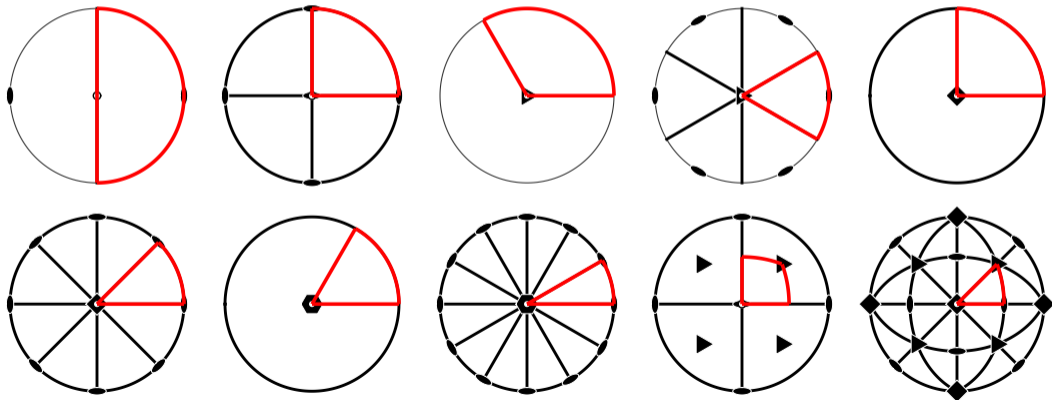
MTEX

```
% define some crystal direction  
cs = crystalSymmetry( '321' , [4.9 4.9 5.4] , 'mineral' , 'quartz' )  
h = Miller ( {1,0,-1,0} , {0,0,0,1} , cs , 'UTW' )  
  
% define some crystal plane  
h = Miller ( {1,0,-1,0} , {0,0,0,1} , cs , 'HKIL' )  
  
% find all symmetrically equivalent  
hSym = h.symmetrise  
unique(h.symmetrise , 'noSymmetry' )  
  
h.multiplicity  
  
angle(hSym(1) , hSym(2))  
  
angle(hSym(1) , hSym(2) , 'noSymmetry' )
```

The Fundamental Sector

Definition

The **fundamental sector** is a spherical region which contains from each class of symmetrically equivalent vectors exactly one.



MTEX

```
% the fundamental sector for a given symmetry
```

```
sR = cs.fundamentalSector
```

```
plot(sR)
```

```
% check whether we are inside the fundamental region
```

```
sR.checkInside(h)
```

```
% define some crystal plane
```

```
h = Miller({1,0,-1,0},{0,0,0,1},cs,'HKIL')
```

```
h = h.project2FundamentalRegion
```

```
% generate vectors within a spherical region
```

```
r = vector3d.rand(sR)
```

```
r = equispacedS2Grid(sR)
```